

I.

Warsaw Talk: Sept 16, 2019

”Standard Classical Mechanics Sitting
in Standard Classical GR”

Ted Newman - Univ of Pittsburgh

Comments:

1. A lot of material - much must be left out - even a
few little dishonesties to save time.

2. Material has appeared in 4-5 published papers
Feb 2019, General Relativity and Gravitation
Aug 2018, Living Reviews

3. We are dealing with **Einstein-Maxwell theory**

**No strings attached,
No non-Commuting VARIABLES,
No higher dimensions**

4. Most of the discussion/action takes place in
the far field region -
In neighborhood of Future Null Infinity

II. I. PRELIMINARIES

1. **Minkowski space-time**, a natural set of null surfaces (in neighborhood of null infinity) are the asymptotic **light cones with origins at arbitrary spacetime points, x^a** - they label the surfaces.

ie. a four real parameter set of these asymptotic null surfaces.

The **optical parameters of associated family of null generators (the null geodesics) - vanishing [TWIST] AND [SHEAR].** .

To each one-parameter family of these asymptotic light cones we can construct a one parameter family of **associated null coordinate systems.**

Now Asymptotically flat space-times.

2. The **major development in gravitational radiation theory** (1960s) was Bondi's introduction of one-parameter families of null surfaces (Bondi surfaces) in the neighborhood of future null infinity.

The null generators of the surfaces (the null geodesics) **HAVE VANISHING TWIST** but have **NON-VANISHING asymptotic SHEAR.**

[Aside; This **asymptotic shear**, $\sigma^0(u, \zeta, \bar{\zeta})$, is the **free data** for the description of gravitational radiation.]

3. Transitioning from **Minkowski space to asymptotically flat space** - one could choose to use the null geodesics that are [either] **TWIST-FREE OR SHEAR FREE.**

Bondi (**naturally**) chose the **TWIST FREE path.**

We suggest its preferable to chose the other path, **I.E. use asymptotically SHEAR FREE** null geodesics near null infinity - **RATHER** then **TWIST-FREE.**

This is THE MAJOR INNOVATION here

III.

4. **CLAIM**; choosing the asymptotically **SHEAR-FREE** version leads to both Bondi's results and to a **NEW** set of remarkable results.

Initially the asymptotically shear-free version is beset with problems, – they can be overcome.

5. **Back to asymptotically flat spacetimes**: Working with the asymptotically shear free null geodesics requires some technology — discussed very superficially.

a. The asymptotically shear free geodesics sets are labeled by four complex parameters, z^a ... Defining a complex 4-space i.e., **H-space**. Each H-space point refers to a **SET** (a bundle) of null rays in the physical spacetime.

The imaginary part of z^a is a measure of the **TWIST** - the real part determines the average position of the set. (a generalization of light-cones.)

We have four real parameters to work with at **INFINITY** - as **IN** the Minkowski case

b. A curve in H-space, $z^a = \xi^a(t) = \xi_R^a(t) + i\xi_I^a(t)$ corresponds in the physical spacetime to a null geodesic congruence with **TWIST**. Each congruence can be used to construct an asymptotic coordinate & tetrad system – an **ASYMPTOTICALLY SHEAR FREE** system.

c. A **VERY** special H-space-curve exists: The **COMPLEX CENTER OF MASS** System

IV. **Quick Review** - HERE just to see what we are talking about

II. - **Asymptotic Einstein-Maxwell Eqs**

1. **Spin-coefficient description of asymptotic Weyl and Maxwell Tensor & Definitions** - Here just to be seen

$$\begin{aligned}\Psi_0 &= \Psi_0^0 r^{-5} + O(r^{-6}), \\ \Psi_1 &= \Psi_1^0 r^{-4} + O(r^{-5}), \\ \Psi_2 &= \Psi_2^0 r^{-3} + O(r^{-4}), \\ \Psi_3 &= \Psi_3^0 r^{-2} + O(r^{-3}), \\ \Psi_4 &= \Psi_4^0 r^{-1} + O(r^{-2}).\end{aligned}$$

$$\begin{aligned}\phi_0 &= \phi_0^0 r^{-3} + O(r^{-4}), \\ \phi_1 &= \phi_1^0 r^{-2} + O(r^{-3}), \\ \phi_2 &= \phi_2^0 r^{-1} + O(r^{-2}),\end{aligned}$$

with

$$\begin{aligned}\Psi_n^0 &= \Psi_n^0(u, \zeta, \bar{\zeta}), \\ \phi_n^0 &= \phi_n^0(u, \zeta, \bar{\zeta}).\end{aligned}$$

The remaining (non-radial) **Bianchi Identities and Maxwell equations** yield the evolution equations:

$$\dot{\Psi}_2^0 = -\Psi_3^0 + \sigma^0 \Psi_4^0 + k \phi_2^0 \bar{\phi}_2^0, \quad (1)$$

$$\dot{\Psi}_1^0 = -\Psi_2^0 + 2\sigma^0 \Psi_3^0 + 2k \phi_1^0 \bar{\phi}_2^0, \quad (2)$$

$$\dot{\Psi}_0^0 = -\Psi_1^0 + 3\sigma^0 \Psi_2^0 + 3k \phi_0^0 \bar{\phi}_2^0, \quad (3)$$

$$k = 2Gc^{-4}, \quad (4)$$

--

$$\dot{\phi}_1^0 = -\phi_2^0, \quad (5)$$

$$\dot{\phi}_0^0 = -\phi_1^0 + \sigma^0 \phi_2^0. \quad (6)$$

Overdot denotes u-derivative.

V.

DEFINITIONS OF PHYSICAL VARIABLES

- a. (a little lie) **Def 1, Bondi-Sachs mass and linear-Momentum** l=0&1
 harmonics of Ψ_2^0 **(Classical)**
 All constants taken as =1, i.e. c=h=G=k=1

$$\Psi_2^0 = M + P^i Y_{1i} \quad (7)$$

- b. **Def 2 Complex Mass Dipole: NEW**
 Mass Dipole plus i Angular Momentum -

$$(D_{(complex)}^i = D_{(mass)}^i + ic^{-1}J^i),$$

$l = 1$ harmonic of Ψ_1^0 ;

$$\Psi_1^0 = (D_{(mass)}^i + ic^{-1}J^i)Y_{1i}^1 + \dots \quad (8)$$

- Def 3** Complex E&M dipole, (electric and i magnetic dipoles,
 $D_{complex}^i = (D_{Elec}^i + iD_{Mag}^i)$
 the $l = 1$ harmonic component of ϕ_0^0 . **(STANDARD)**

$$\phi_0^0 = 2(D_{Elec}^i + iD_{Mag}^i)Y_{1i}^1. \quad (9)$$

VI.

III. MAJOR STEP

Choose a complex world in H-Space, $z_{CofM}^a = \xi^a(t) = \xi_R^a(t) + i\xi_I^a(t)$ with its associated ASYMPTOTICALLY SHEAR FREE system so that **COMPLEX MASS DIPOLE VANISHES - DEFINITION**

$$D_{(complex)}^i = (D_{(mass)}^i + ic^{-1}J^i) = 0 \quad (10)$$

This H-Space curve (and its physical space uniquely associated null geodesic congruence) **define the COMPLEX CENTER OF MASS** .

and a unique **ASYMPTOTICALLY SHEAR FREE** system - **CENTER OF MASS SYSTEM.**

the

VII.

IV. **RESULTS - & *no more definitions***

BY transforming from the **CENTER OF MASS SYSTEM** back to a **Bondi system** we obtain a series surprising and remarkable results.

Result #1 Expressions for Mass Dipole and Spin and Orbital Angular Momentum - **NOT DEFINITIONS** but derived

$$D_{(mass)}^i = M_B \xi_R^i - c^{-1} P_B^k \xi_I^j \epsilon_{jki} + \dots, \quad (11)$$

$$J^i = c M_B \xi_I^i + P_B^k \xi_R^j \epsilon_{jki} + \dots \quad (12)$$

or

$$\vec{D}_{(mass)} = M_B \vec{r} + c^{-2} M_B^{-1} \vec{P}_B x \vec{S}, \quad (13)$$

$$\vec{r} = \xi_R^i = (\xi_R^1, \xi_R^2, \xi_R^3), \quad (14)$$

$$\vec{S} = c M_B \xi_I^j = c M_B (\xi_I^1, \xi_I^2, \xi_I^3), \quad (15)$$

$$\vec{J} = \vec{S} + \vec{r} x \vec{P}. \quad (16)$$

Result #2 From Bianchi Identities **the Kinematic Linear Momentum**

$$P_B^i = M_B \xi_R^{i'} - \frac{2q^2}{3c^3} \xi_R^{i''} \quad (17)$$

Result #3 in Bianchi Identities - **Angular Momentum Conservation.**

$$J^{i'j} = \frac{2q^2}{3c^3} (\xi_R^{j'} \xi_R^{k''} + \xi_I^{k'} \xi_I^{j''}) \epsilon_{kji}, \quad (18)$$

Exactly the same as L & L plus spin loss- no derivation -
JUST sitting in the BI.

VIII.

Result: V - Energy loss

$$M'_B = -\frac{G}{5c^7}(Q_{Mass}^{jk''' } Q_{Mass}^{jk''' } + Q_{Spin}^{jk''' } Q_{Spin}^{jk''' }) - \frac{4q^2}{3c^5}(\xi_R^{i'' } \xi_R^{i'' } + \xi_I^{i'' } \xi_I^{i'' }) - \frac{4}{45c^7}(\dots) \quad (19)$$

Bondi Energy loss & E&M dipole & quadrupole energy loss

Result: 5 - Newton's 2nd Law

$$P_B^{i'} = F_{recoil}^i \quad (20)$$

F_{recoil}^i has many non-linear radiation terms – time derivatives of the gravitational quadrupole and the E&M dipole and quadrupole moments.

Substitute **momentum expression**, Eq. ($P_B^i = M_B \xi_R^{i'} - \frac{2q^2}{3c^3} \xi_R^{i''}$), into **Momentum lose** Eq. ($P_B^{i'} = F_{recoil}^i$) leading to Newton's second law – with Rocket Force and Radiation Reaction Force.

******Result: 6 - Rocket Force and Radiation Reaction Force******

$$M_B \xi_R^{i'''} = F^i \equiv M'_B \xi_R^{i'} + \frac{2q^2}{3c^3} \xi_R^{i'''} + F_{recoil}^i. \quad (21)$$

WE believe: remarkable - exact Abraham-Lorentz-Dirac radiation reaction force - **NO mass renormalization. No derivation - just sitting there to be observed in the $l = 1$ part of a BI.**

IX.

Conclusions:

.
These results of C.M. and E&M theory are clearly sitting in GR.

.
It is not at all clear what are the implications or even the meaning?????

.
Do they fit into a quantum theory? A Schrodinger Eq????
If so HOW?

.
What happens with the runaway behavior resulting from the radiation reaction term in the EQS of motion - is there a term that suppresses the runaway behavior. ?????

.
Is it possible to get two body eqs of motion from this type of analysis???? or other classical results????