Conformally isometric embeddings and Hawking temperature

Maciej Dunajski

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- Maciej Dunajski, Paul Tod (2019) Conformally isometric embeddings and Hawking temperature, arXiv: 1812.05468, CQG 2019.
- Maciej Dunajski, Paul Tod (2019) Conformal and isometric embeddings of Gravitational Instantons, Preprint.



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Conformal embeddings

MANIFOLDS THROUGHOUT THE CENTURIES

• 19th century. Surfaces



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• 20th century. Atlases



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• The Whitney embedding theorem: any n-dimensional manifold can be embedded in \mathbb{R}^N as a surface, where N is at most 2n.

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 - Fubini-Study metric on CP²: embedding class still not known (neither local not global!). At least 3, at most 4.

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• An immersion $\iota : (M,g) \to \mathbb{R}^N$ such that $\iota^*(\eta) = \Omega^2 g$ for some $\Omega : M \to \mathbb{R}^+$, and $\iota(M) \subset \mathbb{R}^N$ is diffeomorphic to M.

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- Naive counting: N embedding functions X^1, \ldots, X^N of local coordinates x^1, \ldots, x^n such that $g = g_{ab}(x)dx^a dx^b$.

$$\eta_{\alpha\beta}\frac{\partial X^{\alpha}}{\partial x^{a}}\frac{\partial X^{\beta}}{\partial x^{b}} = \Omega^{2}g_{ab}, \quad \alpha, \beta = 1, \dots, N, \quad a, b = 1, \dots, n.$$

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This talk:

- **(**) Global conformal embedding of the Schwarzchild metric.
- Obstructions to conformal embeddings of class 1
- Hawking and Unruh temperatures.

• Given an Einstein Lorentzian four-manifold (M,g), seek an isometric embedding of $\hat{g} = \Omega^2 g$ into \mathbb{R}^5 , with second fundamental form

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• Conformal rescallings and spinors: $\hat{C}^{d}_{abc} = C^{d}_{abc}$, $\hat{\sigma}_{ab} = \Omega \sigma_{ab}$ $C_{abcd} = \psi_{ABCD} \epsilon_{A'B'} \epsilon_{C'D'} + \psi_{A'B'C'D'} \epsilon_{AB} \epsilon_{CD}$.

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• Theorem 2. A local conformal embedding ι of Theorem 1, such that rank (K_{ab}) is maximal at some $p \in M$, is rigid in a neighbourhood of p up to conformal transformations of $\mathbb{R}^{r,s}$, r + s = 5.

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LETTER TO THE EDITOR

Twistor equation in a curved spacetime

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Abstract. The twistor equation is studied in a four-real-dimensional spacetime. All the metric tensors which locally admit a solution are found. They either belong to the Fefferman class or are given by the Trautman-Kerr-Schild anzatz by using a non-twisting null conformal Killing vector field in the Minkowski spacetime. The corresponding solutions are derived.

• Algebraic invariants of the Weyl tensor

 $I = \psi_{ABCD} \psi^{ABCD}, \quad J = \psi_{AB}{}^{CD} \psi_{CD}{}^{EF} \psi_{EF}{}^{AB}.$

Algebraically special $J^2 - 6I^3 = 0$. Type 3, or type N: I = J = 0.

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- Corollary: A Riemannian manifold with self-dual Weyl tensor admits a class one conformal embedding iff it is conformally flat.
- The conformal embedding class of CP² is therefore at least two. It is known to be at most three. What is it?

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• $g = V dt^2 - V^{-1} dr^2 - r^2 (d\theta^2 + \sin^2 d\phi^2)$, where V = V(r) has a finite number of simple zeroes $r_0 > r_1 > r_2 \dots$.

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• Find an isometric embedding of $r^{-2}(V^{-1}dr^2 - Vdt^2)$ in AdS_3

$$\frac{dR^2 + dX^2 - dT^2}{R^2}$$

• The unique static, spherically symmetric, global conformal embedding.

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$$h = \exp\Big(\int \frac{V(2V - rV') \pm ar\sqrt{V(4V + 4a^2r^2 - (2V - rV')^2)}}{2rV(a^2r^2 + V)}dr\Big).$$

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- If $V \to 1$ as $r \to \infty$, then $R \sim r$ and $\Omega \sim 1$ as $r \to \infty$.

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•
$$V = 1 - 2m/r$$
, $R(r) = \exp\left(\int \frac{p}{q} dr\right)$, where

$$p = 48m^3 - 16m^2r - r^{3/2}\sqrt{r^3 + 2mr^2 + 4m^2r + 72m^3},$$

$$q = (32m^3 - 16m^2r - r^3)r.$$

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NULL INFINITIES - WHAT HAPPENED TO SCRI?

• Theorem 4. Let $(\mathcal{I}_{\pm})^{Schw}$ and $(\mathcal{I}_{\pm})^5$ be null infinities of the compactified Schwarzschild \overline{M} , and the compactified Minkowski $\overline{\mathbb{R}}^{4,1}$. The conformal embedding extends to a map $\iota : \overline{M} \to \overline{\mathbb{R}}^{4,1}$ s. t. $\iota((\mathcal{I}_{\pm})^{Schw}) = p_{\pm}$ where $p_{-} \in (\mathcal{I}_{-})^5$ and $p_{+} \in (\mathcal{I}_{+})^5$ are points with coordinates $(0, N \subset S^3)$.



• The Hawking effect: a temperature measured by asymptotic observers is $\mathbf{T}_{H}=\kappa/2\pi,$ where

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• The Hawking effect: a temperature measured by asymptotic observers is $\mathbf{T}_{H} = \kappa/2\pi$, where

 $\nabla_a(|K|^2)=-2\kappa K_a,\quad \text{and}\quad g(K,K)=0\quad \text{is a Killing horizon}.$

- The Unruh effect: an observer moving with constant acceleration α in the Minkowski space measures a temperature $\mathbf{T}_U = \alpha/2\pi$.
- Deser and Levin (1999): The Hawking temperature in (M,g) equals the Unruh temperature in an isometric embedding extending through the Killing horizon.
- Holds for the conformal isometric embedding of Schwarzchild in ℝ^{4,1}:
 (M, ĝ) not Einstein, but Hawking effect is kinematical.
 - **2** The surface gravity is conformally invariant as long as Ω and $d\Omega$ are regular on the horizon, Ω is static, and $\Omega \to 1$ when $r \to \infty$.

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 - **③** A trajectory of $K = \partial/\partial t$ in M lifts to a hyperbola in $\mathbb{R}^{4,1}$

$$X_1^2 - X_0^2 = \frac{16m^2h(r)^2}{r^2} \left(1 - \frac{2m}{r}\right) \equiv \alpha^{-2}, \text{ where } r \text{ fixed.}$$

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() Use Tolman's law, take a limit $r \to \infty$.

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Happy Birthday Jurek!

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